

```
[1]: import numpy as np
      from matplotlib import pyplot as plt
```

Course: Computational Electrodynamics and Applications

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0.1 Numerical study of 1D wave equation

- The code required for the study is given below.

```
[2]: def weq_es(u0, um1, c, dt, dx, Nt, Nx):
      '''
      Explicit scheme of second-order (using central differences)
      Given in [GedneyEE624]; solved in vector form
      '''
      cx2 = (c*dt/dx)**2

      u = np.empty((Nt, Nx))
      u[0, :] = u0

      u[1, 1:-1] = -um1[1:-1] + 2*u[0, 1:-1] + \
                   cx2*(u[0, 2:] - 2*u[0, 1:-1] + u[0, :-2])
      u[:, 0] = u[:, -1] = 0

      for n in range(1, Nt-1):
          u[n+1, 1:-1] = -u[n-1, 1:-1] + 2*u[n, 1:-1] + \
                        cx2*(u[n, 2:] - 2*u[n, 1:-1] + u[n, :-2])

      return u
```

```
[3]: def weq_is(u0, um1, c, dt, dx, Nt, Nx):
      '''
      Implicit scheme of second-order (using Newmark scheme)
      Given in [GedneyEE624]; solved in matrix form
      '''
      cx2 = (c*dt/dx)**2

      L = 2*np.identity(Nx) \
          - np.diag(np.ones(Nx-1), -1) \
          - np.diag(np.ones(Nx-1), 1)      #tridiag([-1,2,-1])

      beta = 1/4

      I = np.identity(Nx)
```

```

A = beta*L + (1/cx2)*I
B = (2*beta-1)/2*L + (1/cx2)*I
Ainv = np.linalg.inv(A)

u = np.empty((Nt, Nx))
u[0, :] = u0

u[1, :] = 2*Ainv @ B @ u[0, :] - um1
u[1, 0] = u[1, -1] = 0

for n in range(1, Nt-1):
    u[n+1, :] = 2*Ainv @ B @ u[n, :] - u[n-1, :]
    u[n+1, 0] = u[n+1, -1] = 0

return u

```

- Assuming that

```
[4]: c = 1
```

- Assuming a rectangular pulse as the initial condition.

Not explicitly mentioned, so we consider that $i \in [0, Nx]$, or else an index-0 numbering is used. Under this, pulse starts in the third ($i=2$) space step, where the first ($i=0$) space step should be the boundary. Also, required are the space length. As not an exact number was given we consider that sufficiently big.

```
[5]: L = 10
     Nx = 100
```

```
[6]: def rec_pulse(lo, hi, Nx, ex):
     v = np.zeros(Nx)
     v[ex[0]:ex[1]+1] = 1
     return v

def rec_pulse_mv(v, t):
     hidx = v.argmax()
     ex = [hidx+t, hidx+v[hidx:].argmin()-1+t]
     return rec_pulse(v.min(), v[hidx], np.size(v), ex)

um1 = rec_pulse(0, 1, Nx, [2, 11])
u0 = rec_pulse_mv(um1, 1)

```

- Recording and showing snapshots of the wave.

```
[7]: dx = L/Nx
     n_range = range(20, 60+1, 10)
     Nt = n_range[-1]+1
```

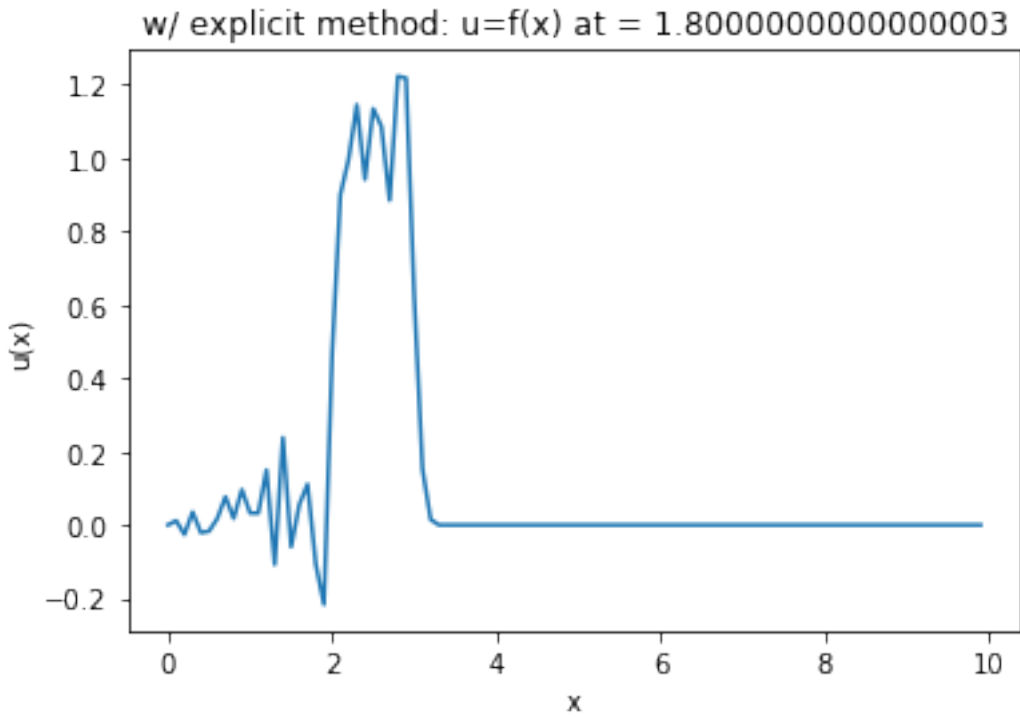
```
[8]: def plot_wrapper(n_range, u0, um1, c, dt, dx, Nt, Nx):
    x = [i*dx for i in range(Nx)]

    u = weq_es(u0, um1, c, dt, dx, Nt, Nx)
    for n in n_range:
        plt.figure()
        plt.plot(x, u[n,:])
        plt.title("w/ explicit method: u=f(x) at = {}".format(n*dt))
        plt.xlabel('x')
        plt.ylabel('u(x)')
        plt.show()

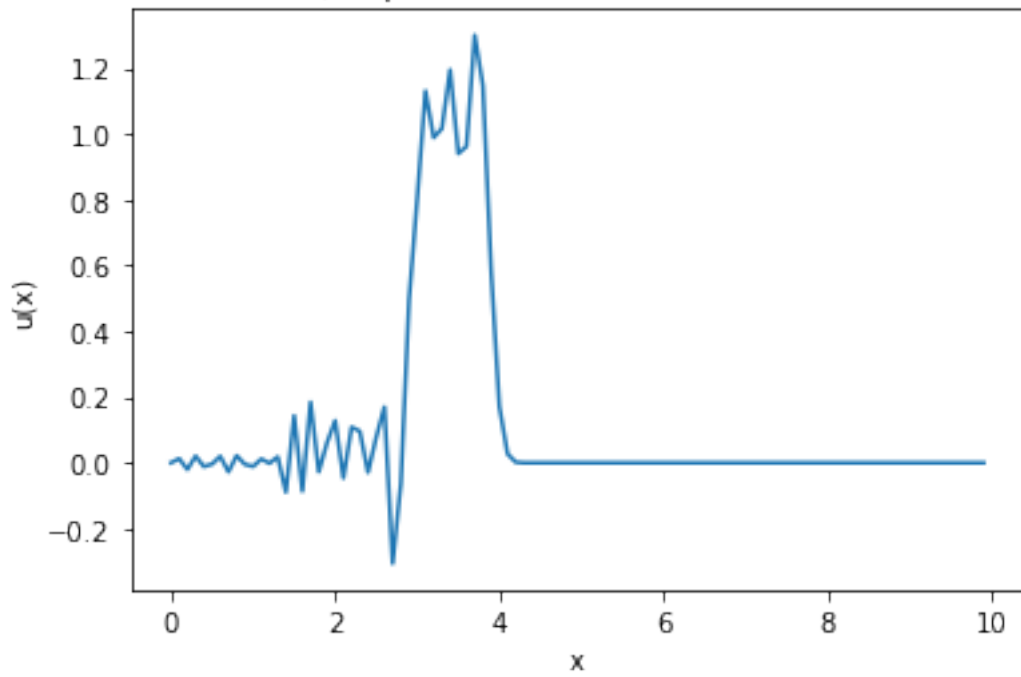
    u = weq_is(u0, um1, c, dt, dx, Nt, Nx)
    for n in n_range:
        plt.figure()
        plt.plot(x, u[n, :])
        plt.title("w/ implicit method: u=f(x) at t = {}".format(n*dt))
        plt.xlabel('x')
        plt.ylabel('u(x)')
        plt.show()
```

For $\Delta t = 0.9\Delta x/c$

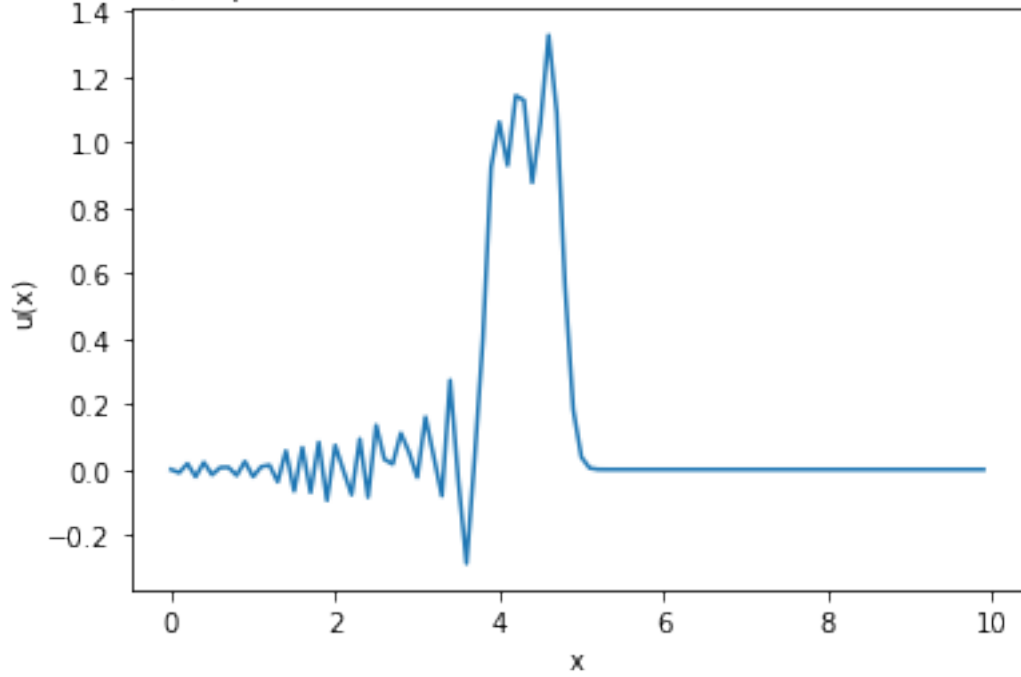
```
[9]: dt = 0.9*dx/c
    plot_wrapper(n_range, u0, um1, c, dt, dx, Nt, Nx)
```



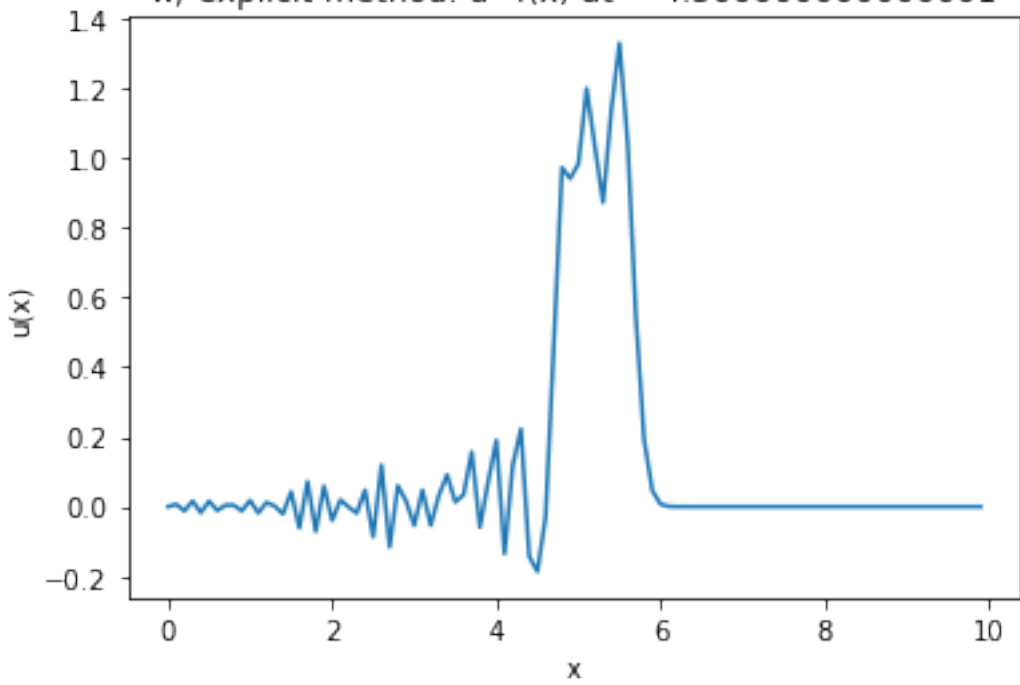
w/ explicit method: $u=f(x)$ at = 2.7



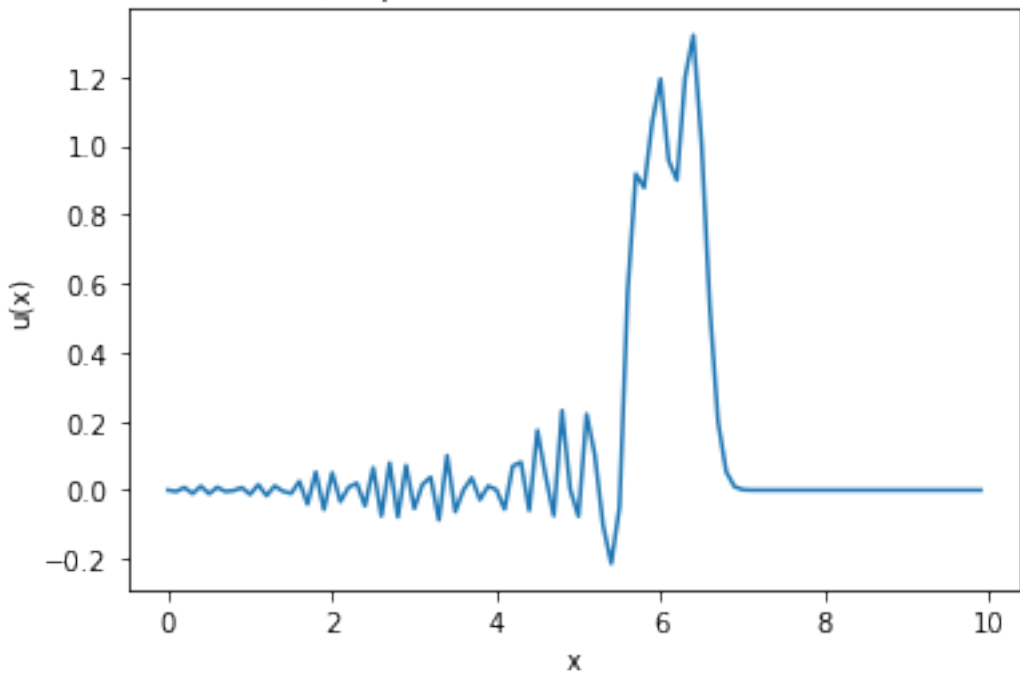
w/ explicit method: $u=f(x)$ at = 3.6000000000000005



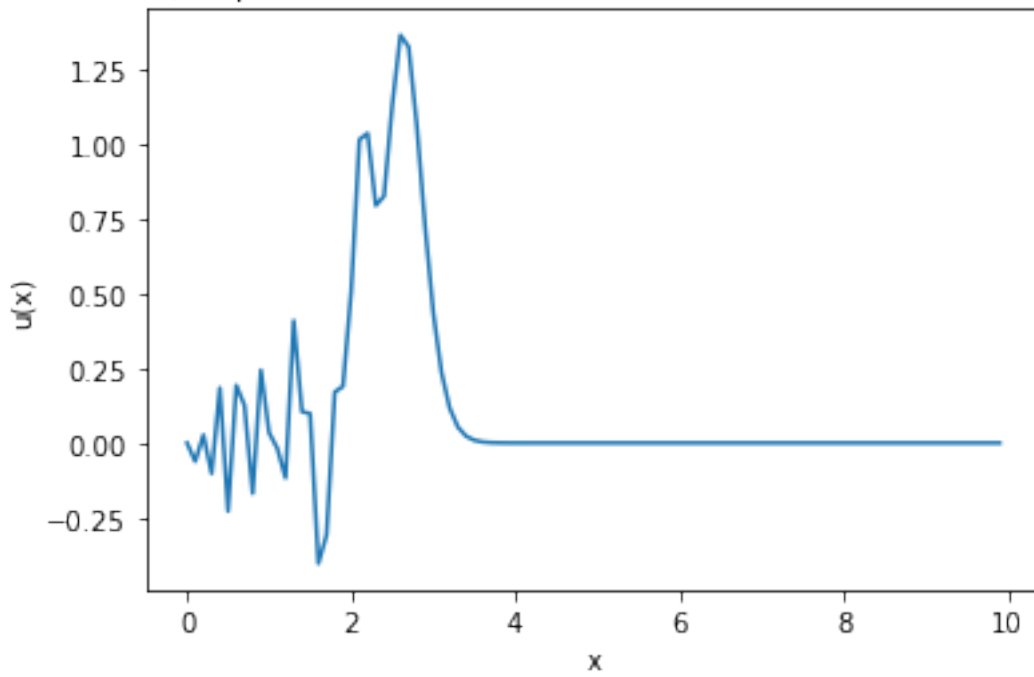
w/ explicit method: $u=f(x)$ at = 4.500000000000001



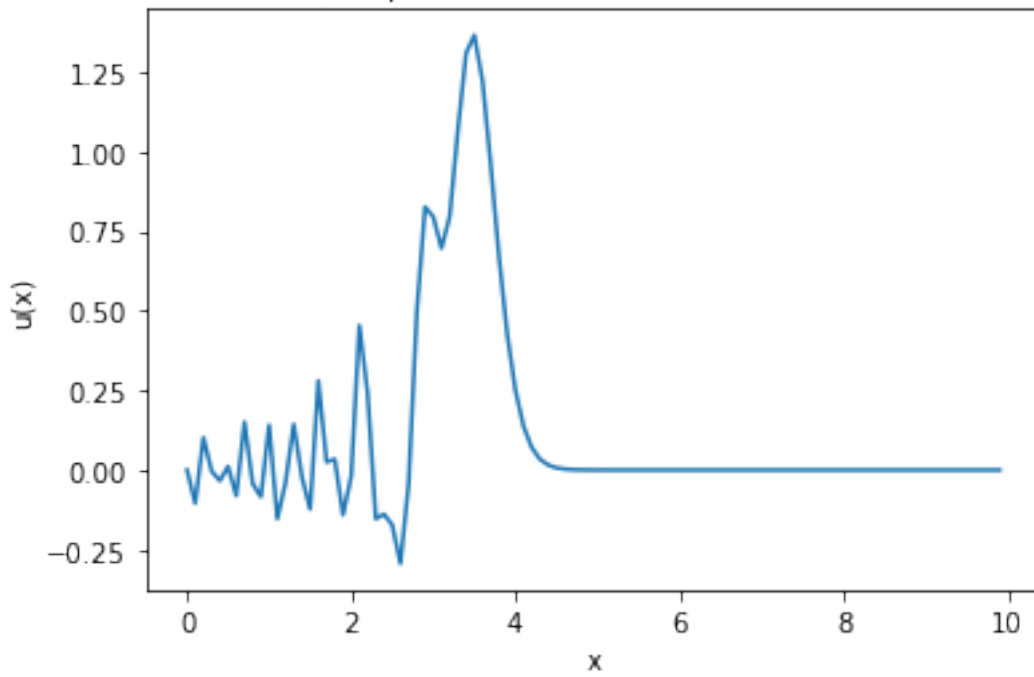
w/ explicit method: $u=f(x)$ at = 5.4



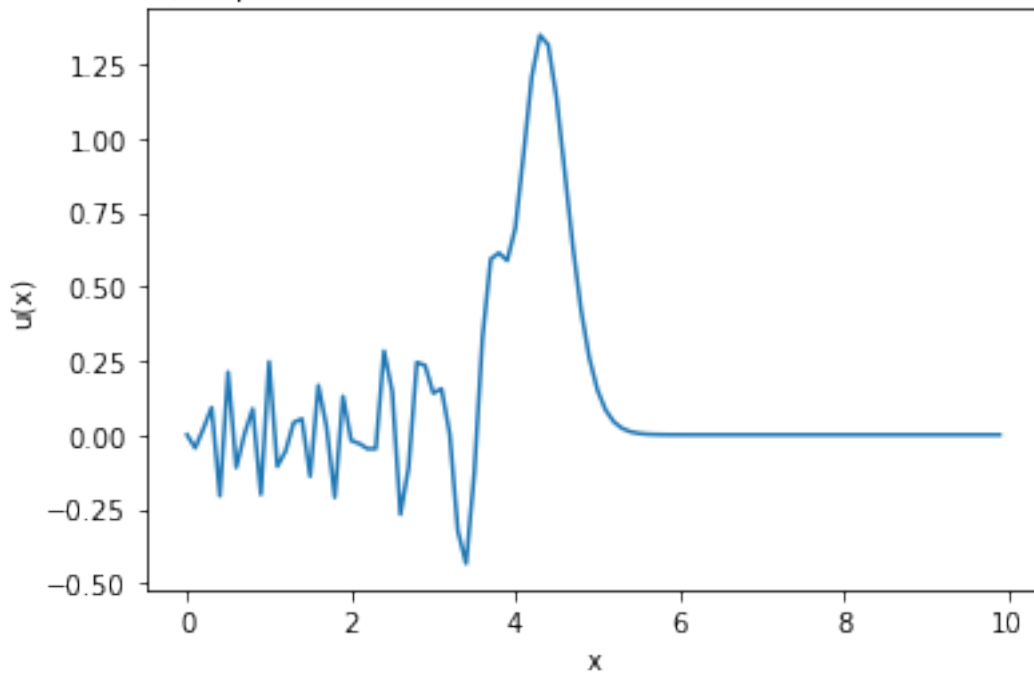
w/ implicit method: $u=f(x)$ at $t = 1.8000000000000003$



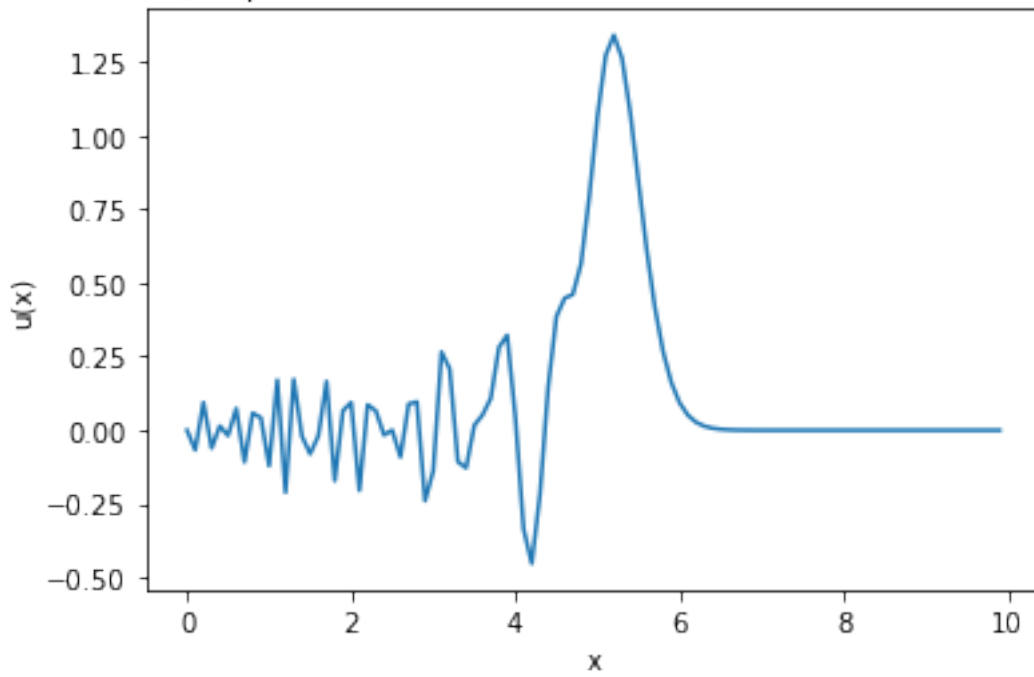
w/ implicit method: $u=f(x)$ at $t = 2.7$

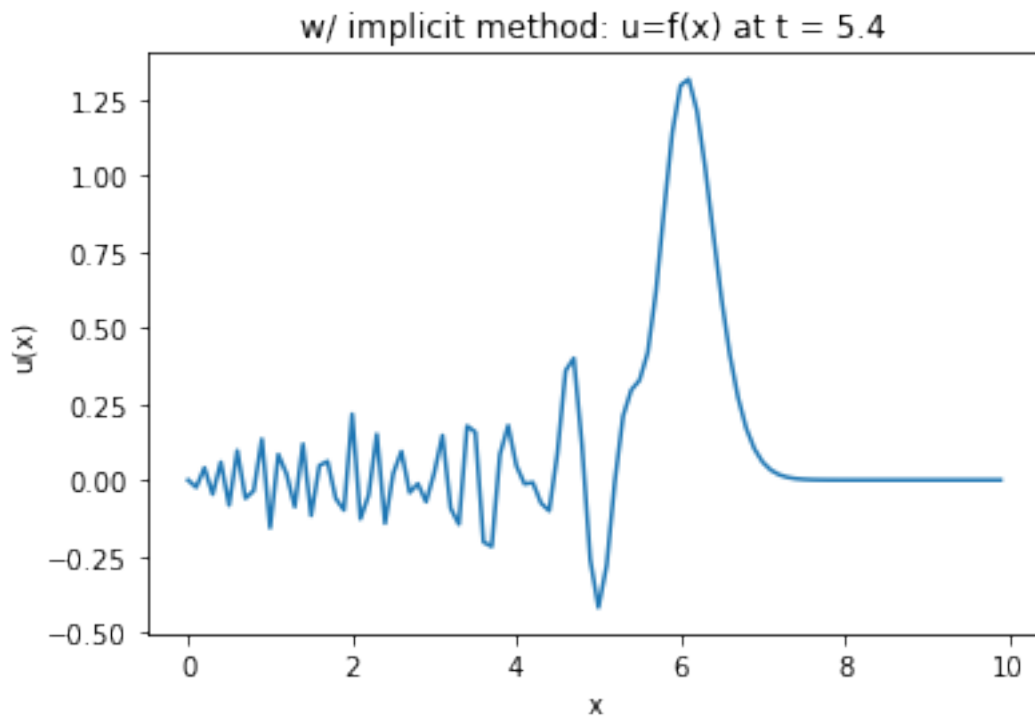


w/ implicit method: $u=f(x)$ at $t = 3.6000000000000005$



w/ implicit method: $u=f(x)$ at $t = 4.5000000000000001$

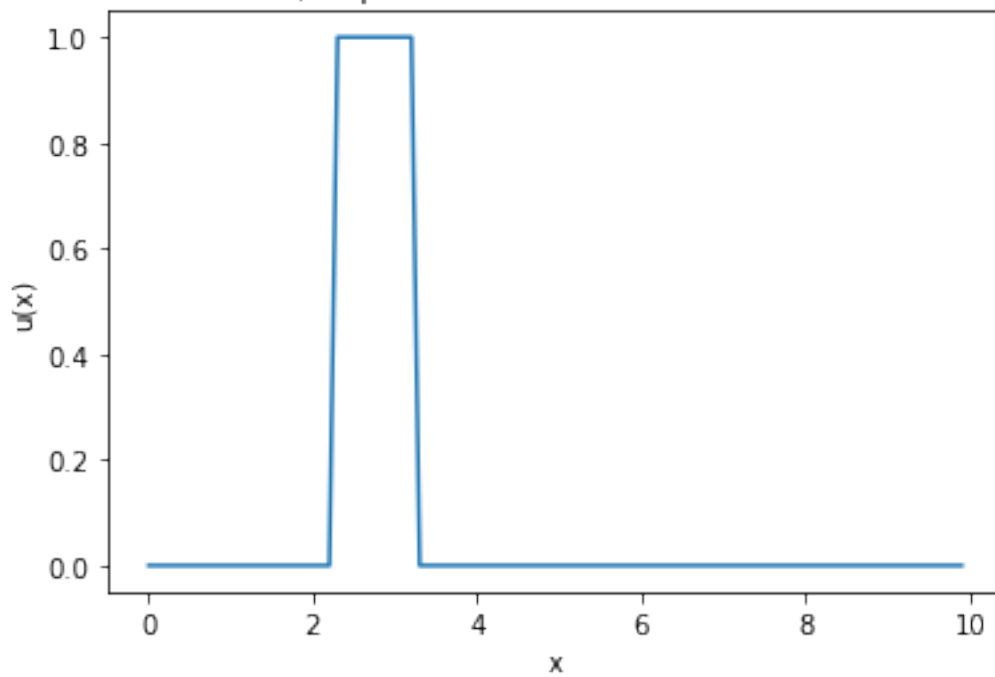




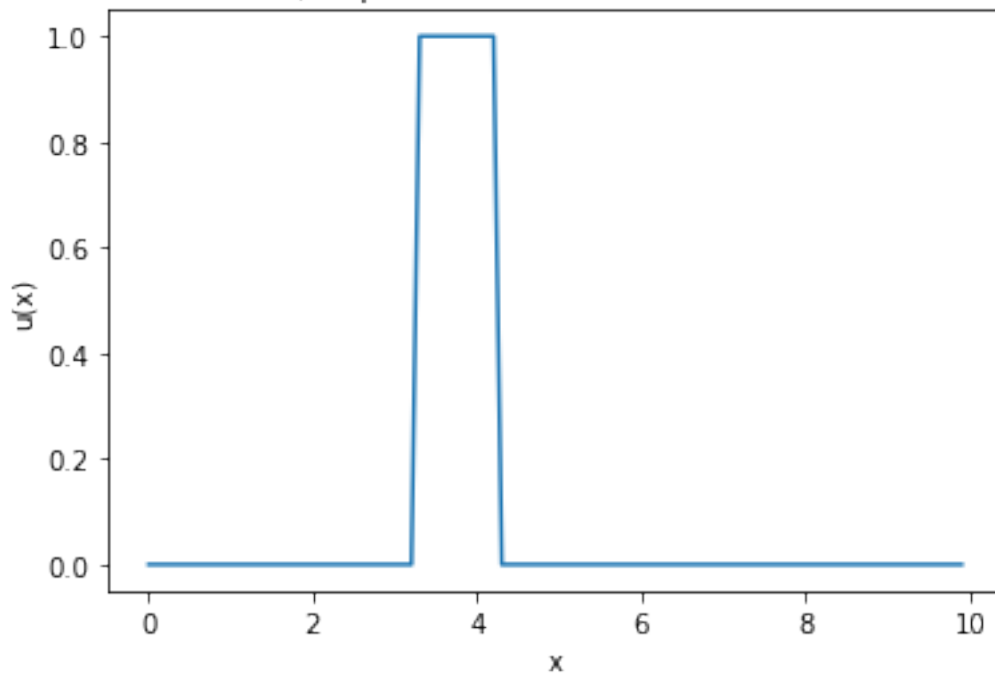
For $\Delta t = \Delta x/c$

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[10]: dt = 1.0*dx/c  
plot_wrapper(n_range, u0, um1, c, dt, dx, Nt, Nx)
```

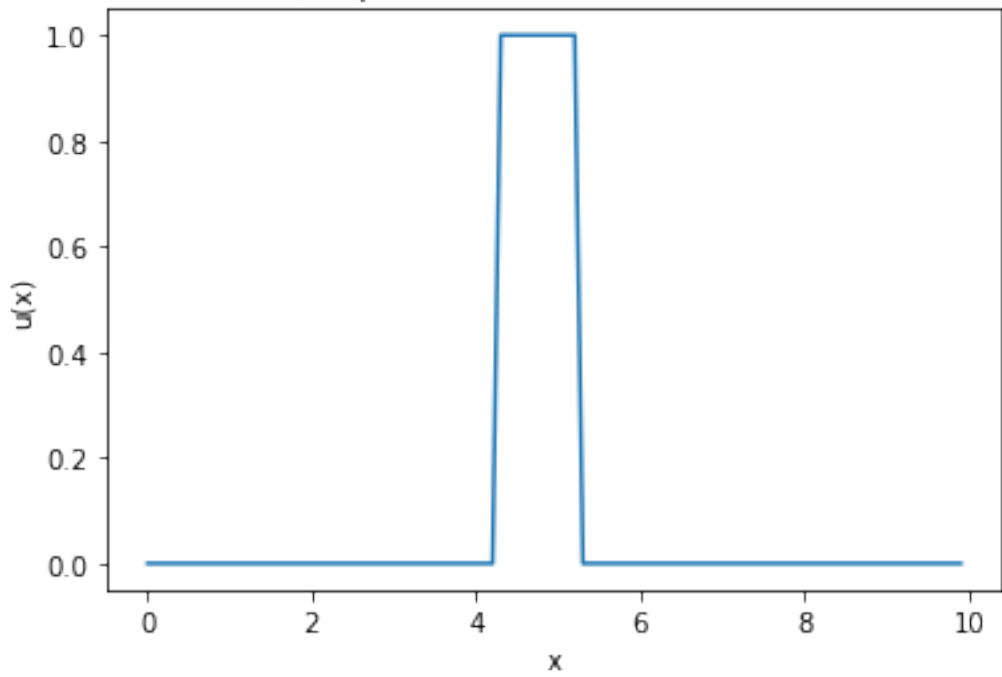

w/ explicit method: $u=f(x)$ at $t = 2.0$



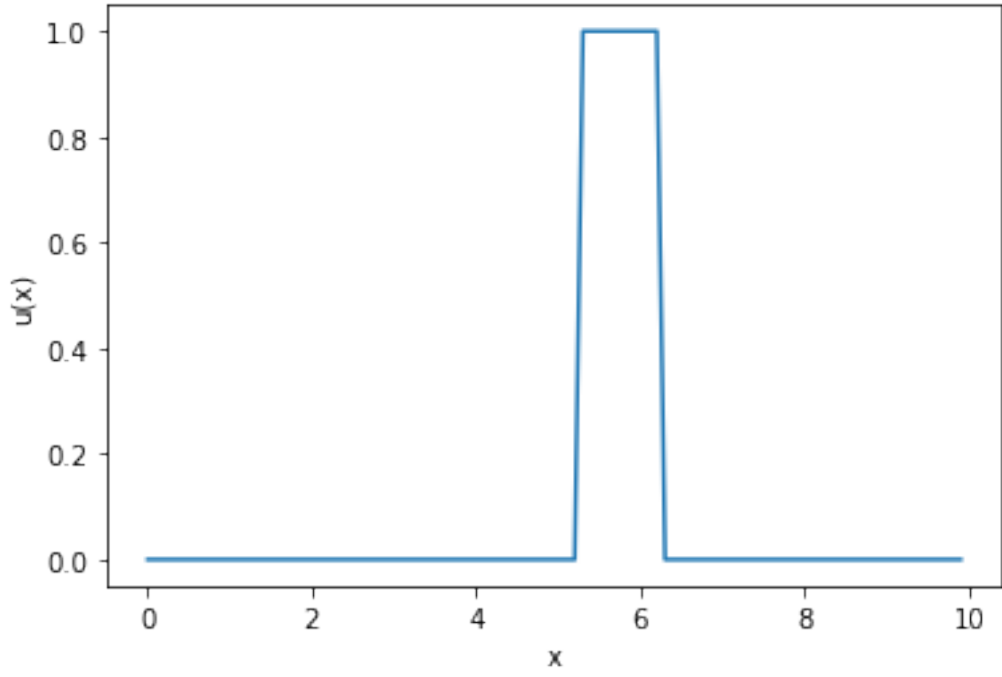
w/ explicit method: $u=f(x)$ at $t = 3.0$



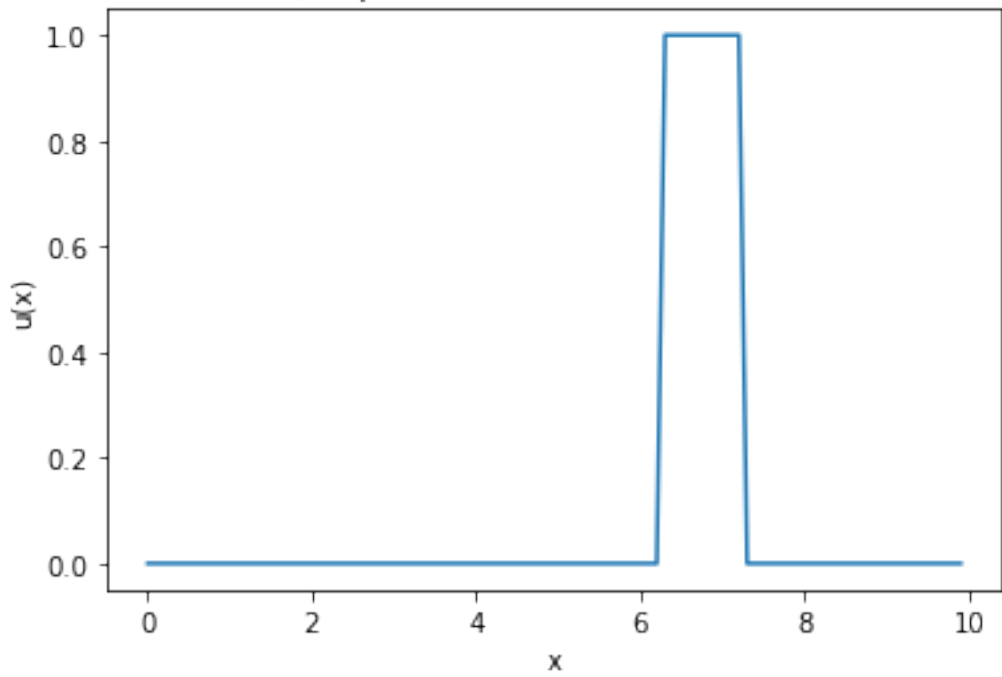
w/ explicit method: $u=f(x)$ at = 4.0



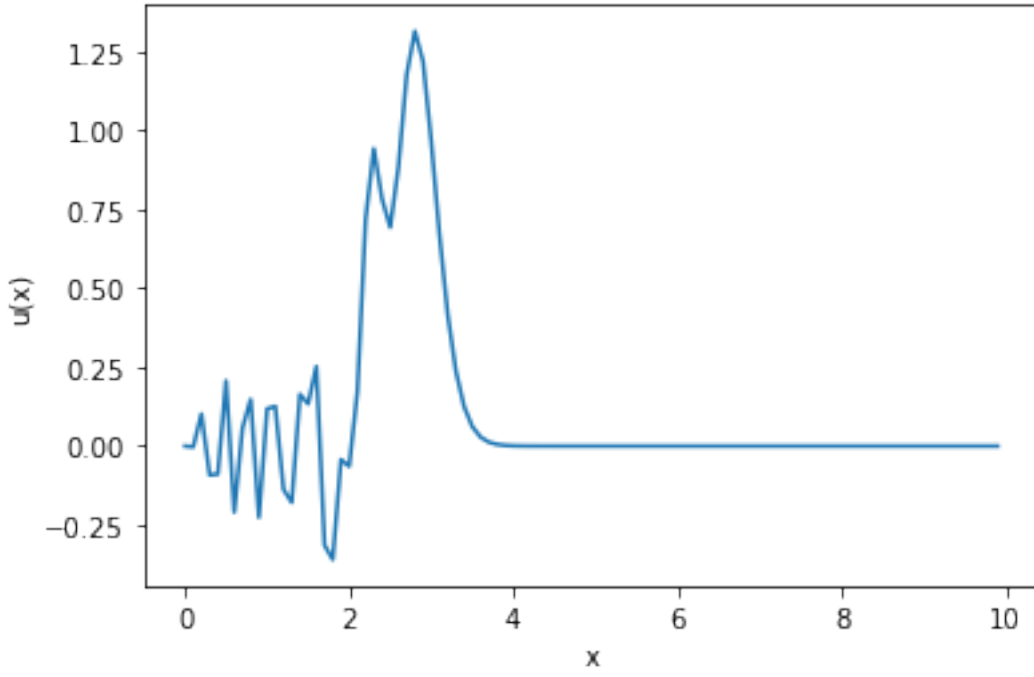
w/ explicit method: $u=f(x)$ at = 5.0

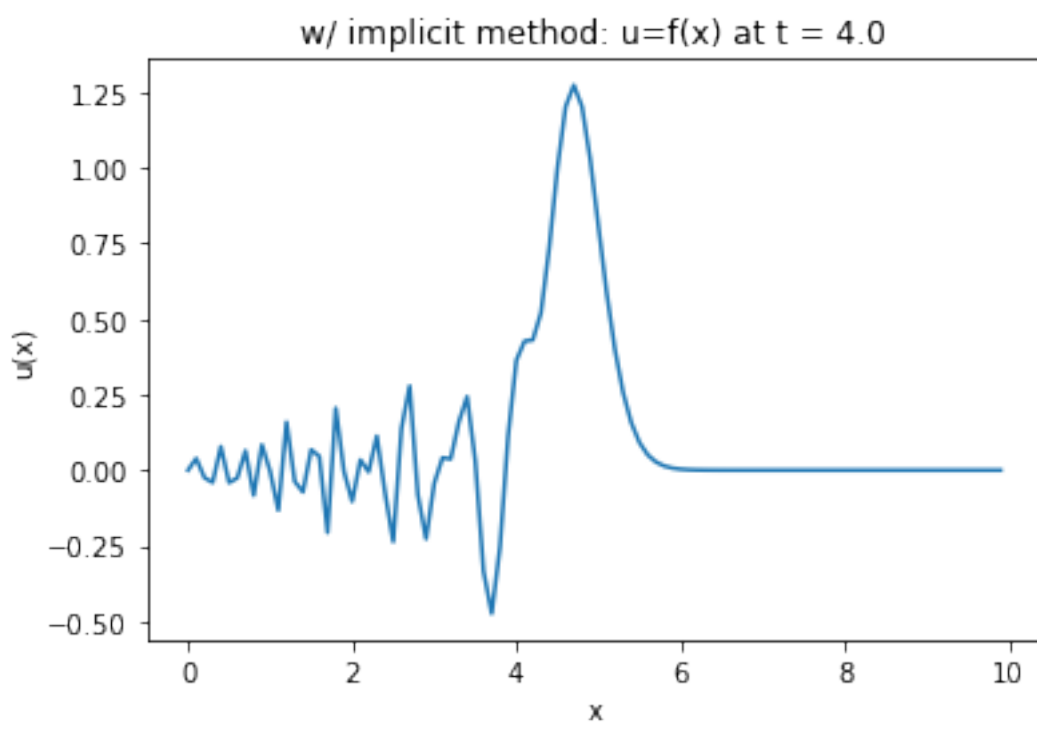
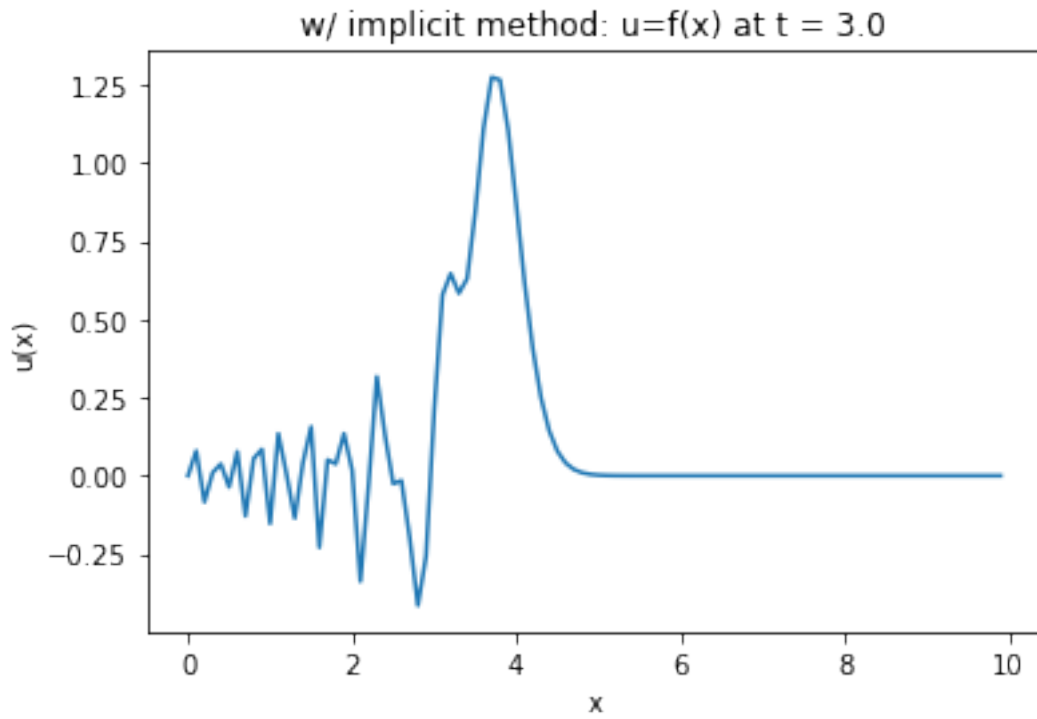


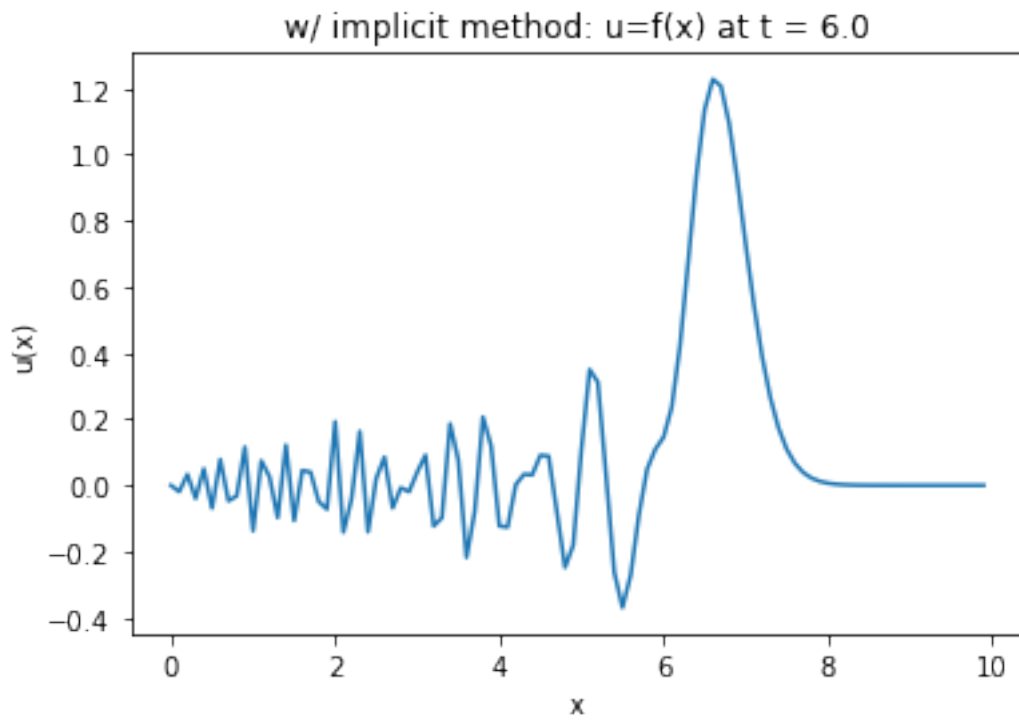
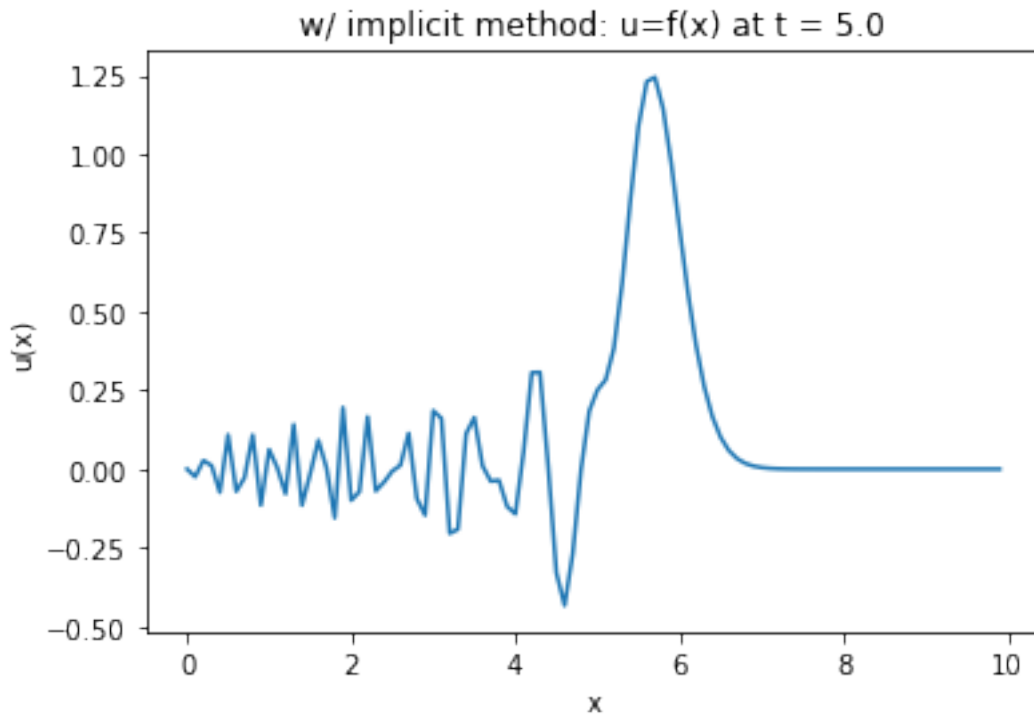
w/ explicit method: $u=f(x)$ at $t = 6.0$



w/ implicit method: $u=f(x)$ at $t = 2.0$

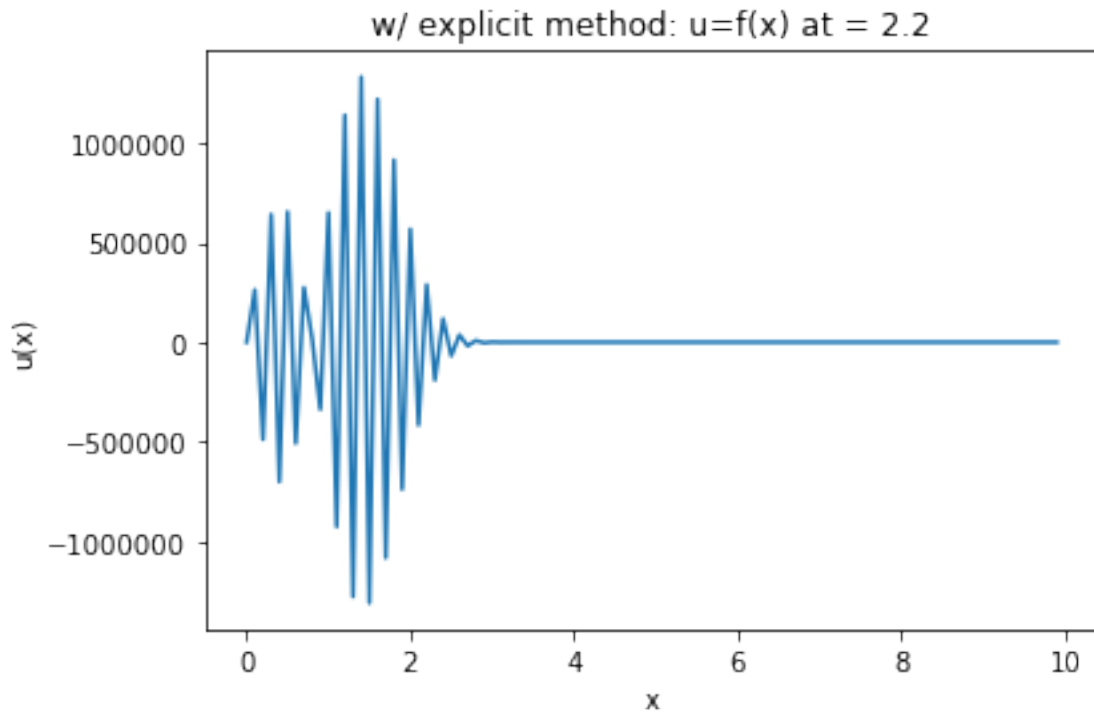




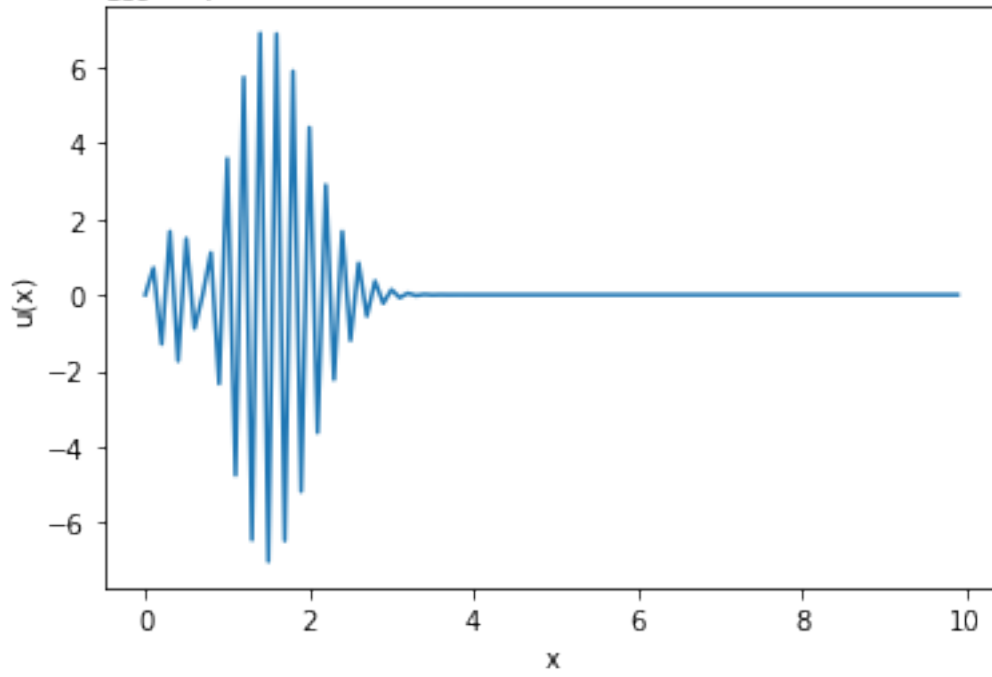


For $\Delta t = 1.1\Delta x/c$

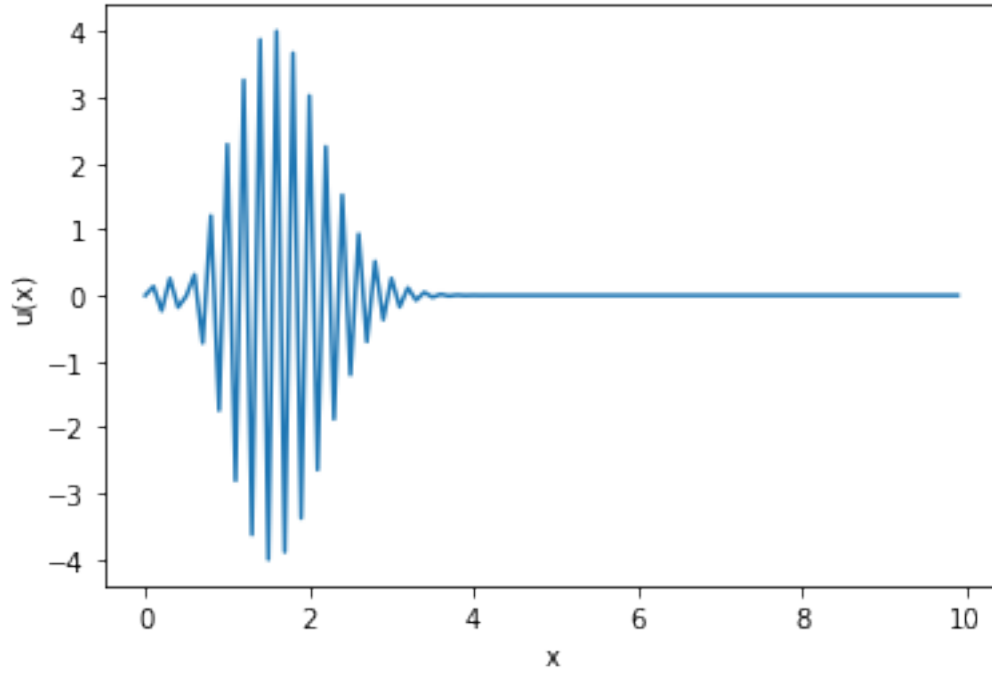
```
[11]: dt = 1.1*dx/c  
plot_wrapper(n_range, u0, um1, c, dt, dx, Nt, Nx)
```



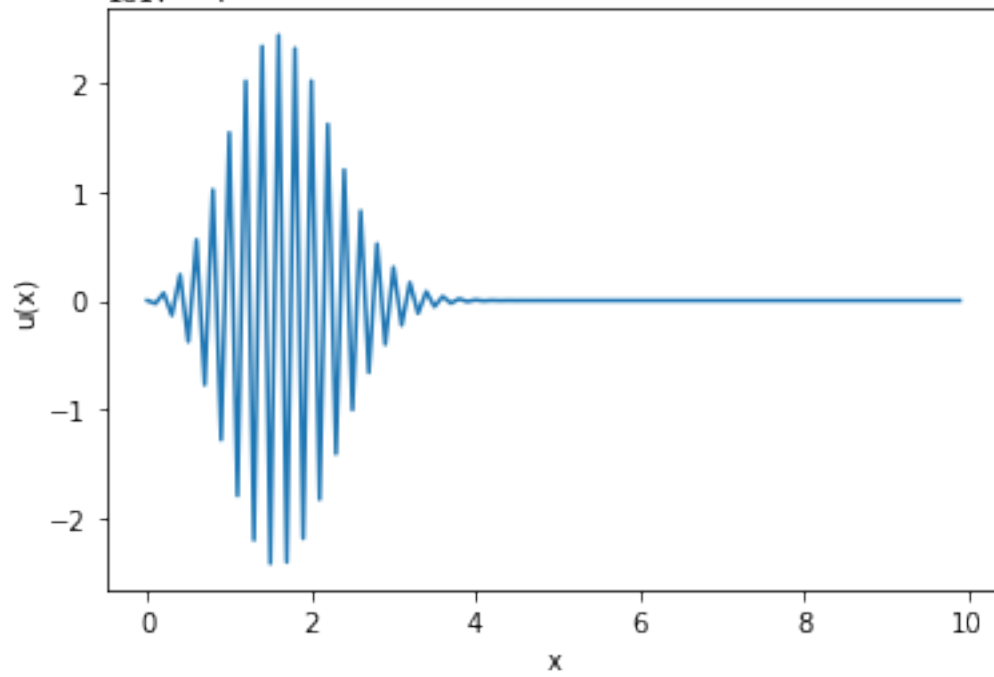
w/ explicit method: $u=f(x)$ at = 3.3000000000000003



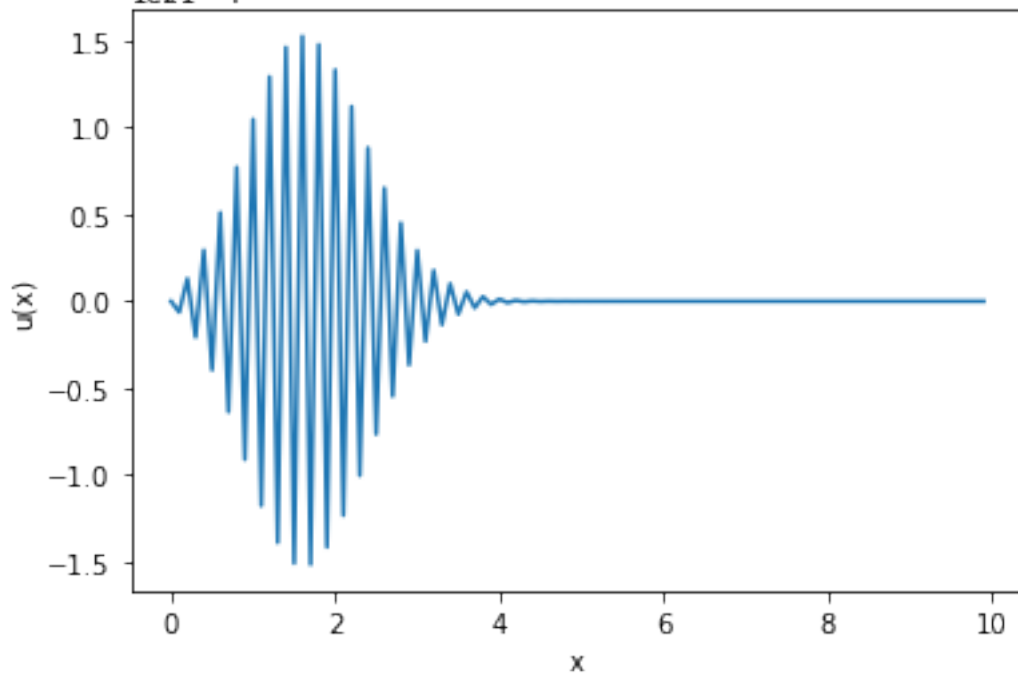
1e13 w/ explicit method: $u=f(x)$ at = 4.4



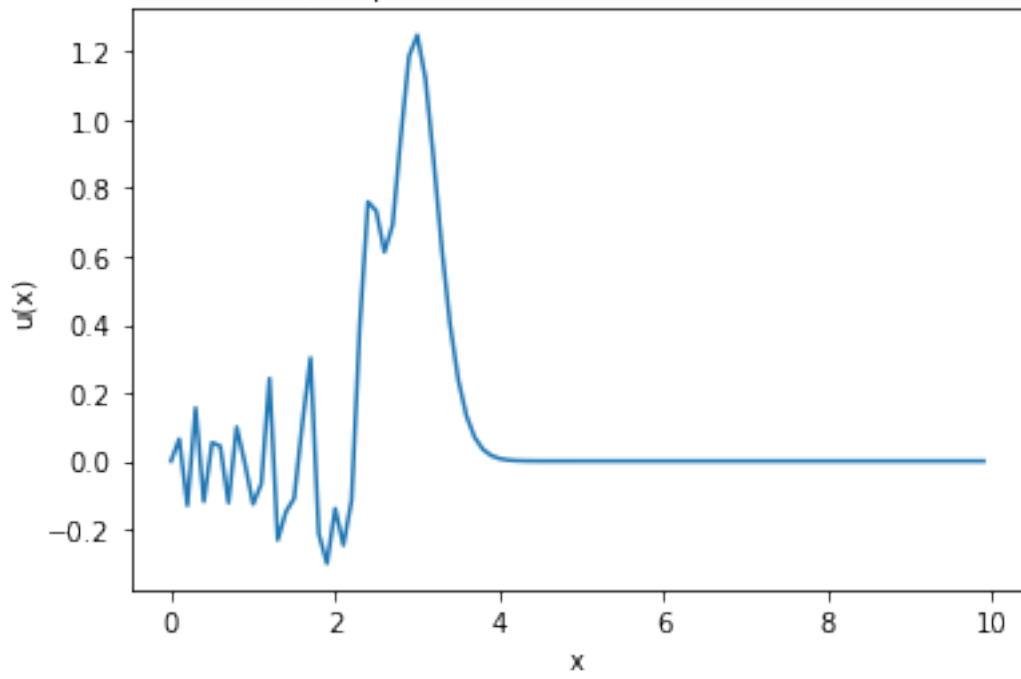
w/ explicit method: $u=f(x)$ at = 5.500000000000001



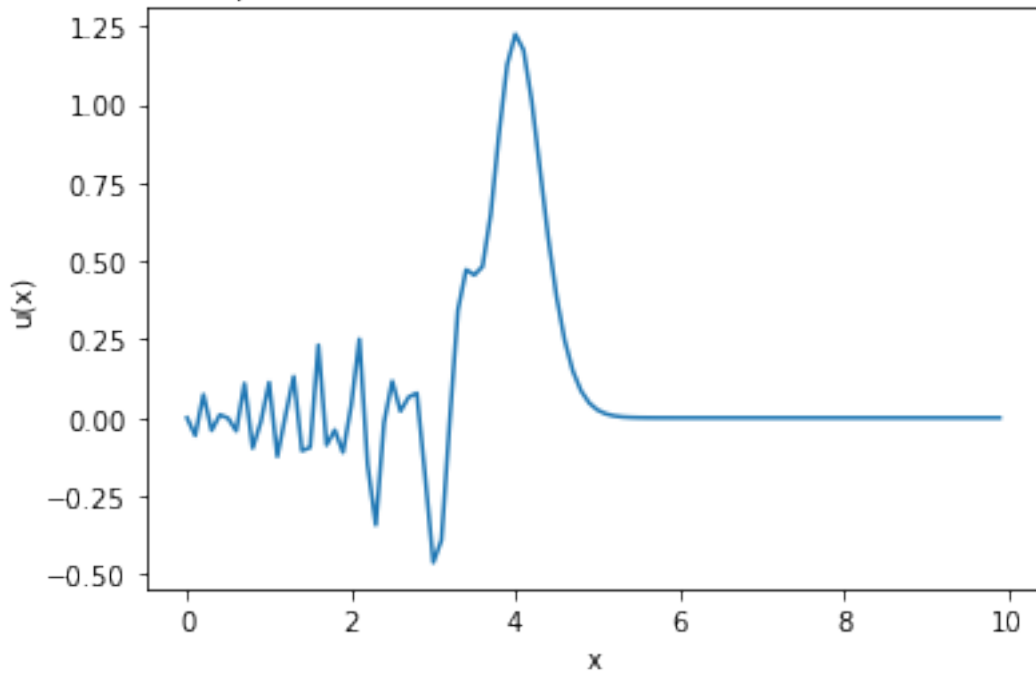
w/ explicit method: $u=f(x)$ at = 6.600000000000005

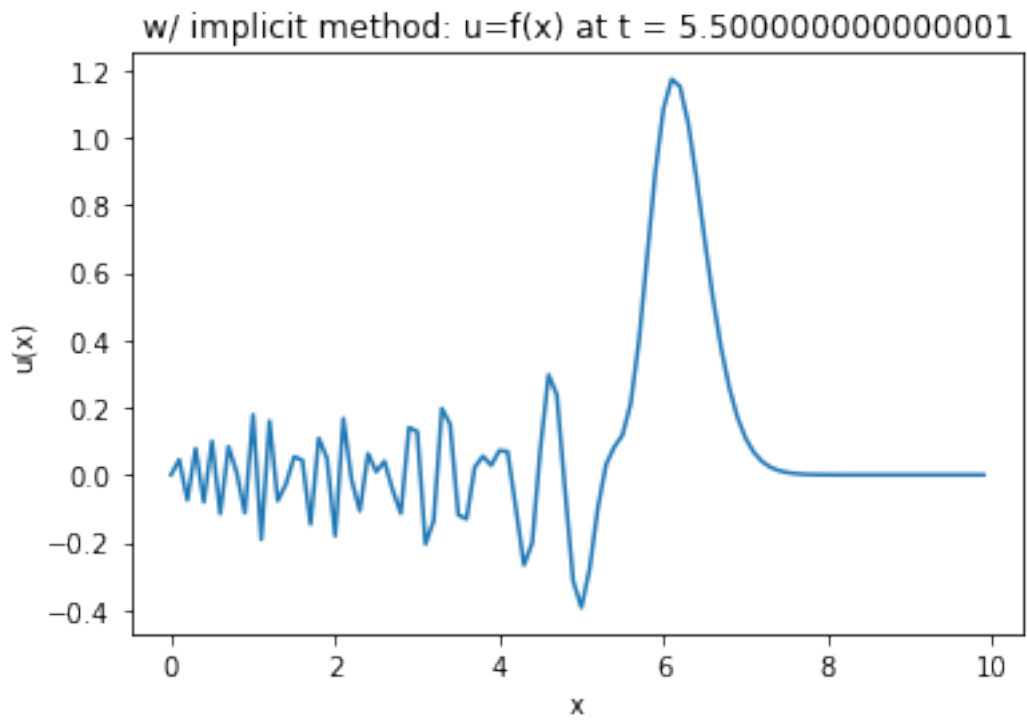
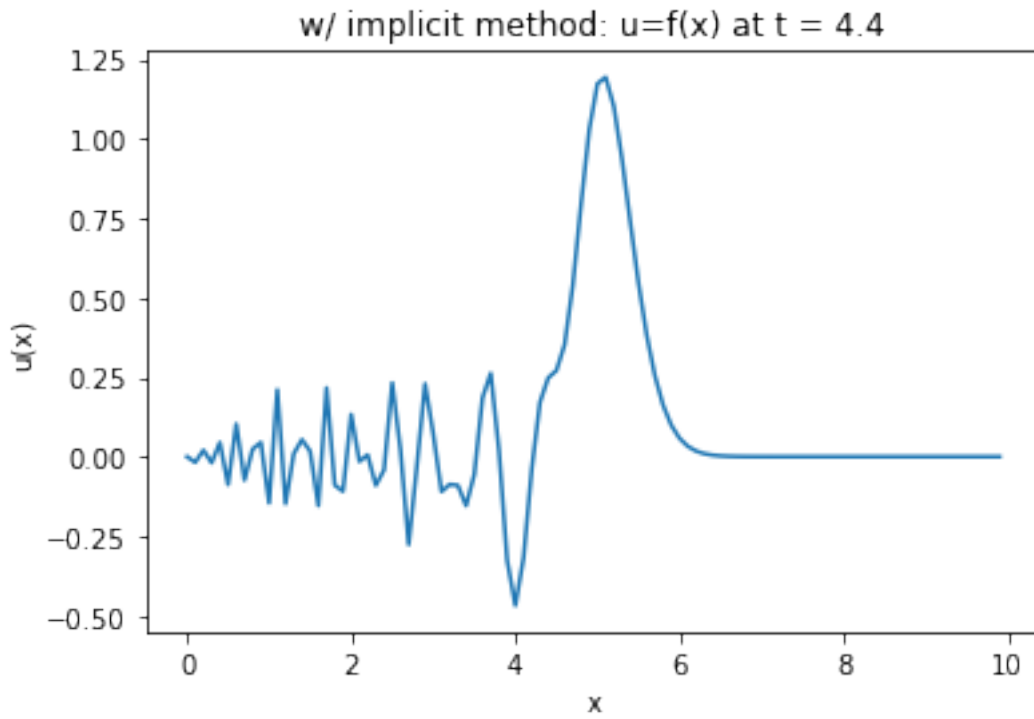


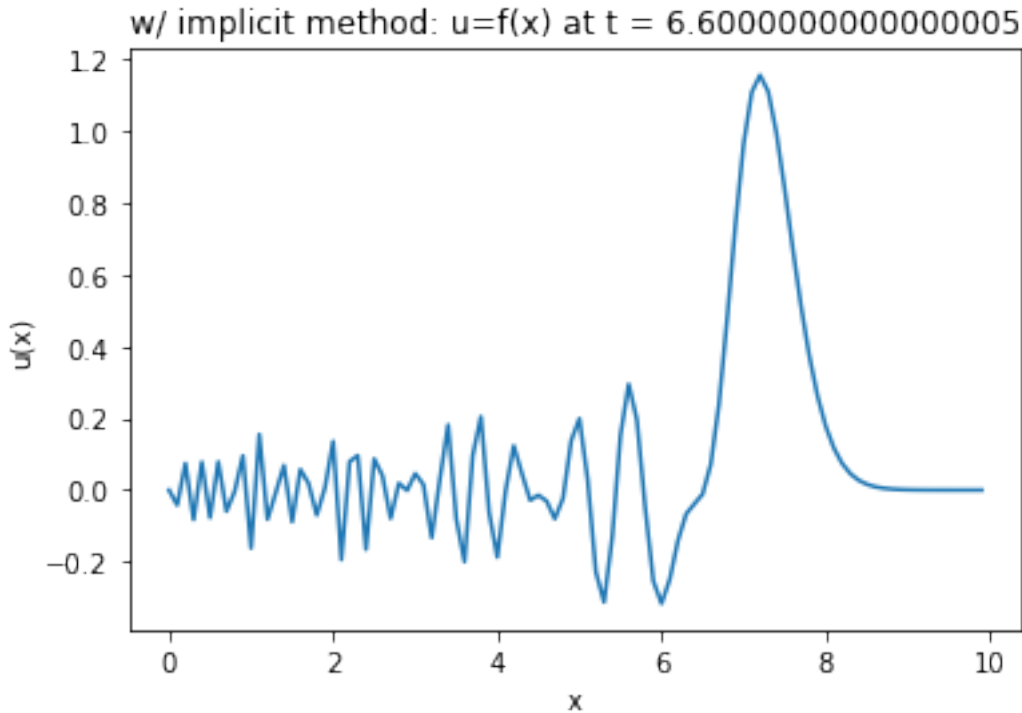
w/ implicit method: $u=f(x)$ at $t = 2.2$



w/ implicit method: $u=f(x)$ at $t = 3.3000000000000003$







- Comments

We tried two 2nd-order methods, an explicit and an implicit, for numerically solving the one-dimensional wave equation, using a time step $dt = a \cdot dx/c$ for various values of a . As the initial conditions represent a rectangular pulse, the integration should result in a moving pulse.

For the explicit scheme, in the first case where $a < 1$ we're seeing ringing due to numerical dispersion. In the second case where $a = 1$ we get an exact rectangular pulse. Dispersion has disappeared as we've taken a time step $= dx/c$, the so-called magic time step.

This magic time step doesn't apply to the implicit scheme. Instead the way to improve the results is taking $dx, dt \rightarrow 0$ (utilizing fine sampling).

In the third and final case where $a > 1$, we notice that the explicit scheme becomes unstable, whereas the implicit scheme continues working. Theoretically this was expected as the stability condition for the explicit scheme is $a = c \cdot dt/dx \leq 1$. In contrast the implicit scheme is absolute stable.

Bibliography

- [GedneyEE624] Gedney Stephen, University of Kentucky, EE624 Notes, Fall 2005, "1D Solution of the Wave Equation".