

Autonomous Mechanical systems of 1 degree of freedom

$$\ddot{x} = f(x, \dot{x}), \quad x \in \mathbb{R}$$

$$\begin{aligned}\dot{x} &= y \\ \dot{y} &= f(x, y)\end{aligned}\quad \begin{pmatrix} \eta' & \dot{x}_1 = x_2 \\ & \dot{x}_2 = f(x_1, x_2) \end{pmatrix}$$

x : position

y : velocity

f : force function

Solution : $x = x(t; c_1, c_2)$, $y = y(t; c_1, c_2)$, $c_i = c_i(x_0, y_0)$

Integral of motion : $F(x, y) = c$, $c = c(x_0, y_0)$

Analytic solutions / integrals ?

Autonomous Mechanical systems of 1 degree of freedom

Example 1.

$$\ddot{x} = x \cdot \dot{x} \quad \dot{\eta} \quad \dot{x} = y, \quad \dot{y} = xy$$

$$\ddot{x} = \frac{d\dot{x}}{dt} = \frac{d\dot{x}}{dx} \frac{dx}{dt} = \frac{dy}{dx} y$$

$$y \frac{dy}{dx} = xy \stackrel{y \neq 0}{\Rightarrow} dy = xdx \stackrel{\int}{\Rightarrow} y = \frac{1}{2}(x^2 + c_1) \Rightarrow F(x, y) = 2y - x^2 = c_1$$

Integral of motion

We consider $c_1 > 0$

$$y = \frac{dx}{dt} = \frac{1}{2}(x^2 + c_1) \Rightarrow \frac{dx}{x^2 + c_1} = 2dt \stackrel{\int}{\Rightarrow} \tan^{-1}(x/\sqrt{c_1}) = 2\sqrt{c_1}t + c_2 \quad \text{or}$$

$$x(t) = \sqrt{c_1} \tan(2\sqrt{c_1}t + c_2) \quad \text{and} \quad y(t) = \frac{dx}{dt} = \frac{2c_1}{\cos^2(2\sqrt{c_1}t + c_2)} \quad (c_1 > 0)$$

$$c_1 = \frac{x_0^2}{\tan^2 c_2}, \quad c_2 = \sin^{-1} \sqrt{\frac{2x_0^2}{y_0}}$$

Autonomous Mechanical systems of 1 degree of freedom

Example 2.

$$\ddot{x} = x(1 + \dot{x})$$

$$\ddot{x} = y dy / dx$$

$$\frac{ydy}{1+y} = xdx \quad \int \quad y - \ln|1+y| = \frac{1}{2}x^2 + c_1$$

$$\Rightarrow y - \ln|1+y| - \frac{1}{2}x^2 = c_1 \quad \text{Integral of motion}$$

$$\Rightarrow \frac{dx}{dt} = y = \dots ?$$

Autonomous Mechanical systems of 1 degree of freedom

$$\ddot{x} = f(x, \dot{x}), \quad \text{or} \quad \begin{aligned} \dot{x} &= y \\ \dot{y} &= f(x, y) \end{aligned}$$

General features

- *Dimensions* : 2 (dynamical variable x and y)
- *Phase Space* : The plane position-velocity (x, y)
- *Autonomous system* : phase space trajectories do not intersect on the plane (x, y)
- *Area Preservation* $\operatorname{div} \vec{f} = \frac{\partial f_1}{\partial x_1} + \frac{\partial f_2}{\partial x_2} = \frac{\partial y}{\partial x} + \frac{\partial f(x, y)}{\partial y} = \frac{\partial f(x, y)}{\partial y}$

$$\operatorname{div} \vec{f} = 0 \iff \frac{\partial f}{\partial y} = 0 \quad \text{conservative systems}$$

Part A. Conservative Systems

$$\ddot{x} = f(x) \quad \text{or} \quad \dot{x} = y, \quad \dot{y} = f(x)$$

- The **energy integral** $E = \frac{1}{2} y^2 + U(x) = E_0 = \sigma \tau \alpha \theta.$, $E_0 = E(x_0, y_0)$
where $U = U(x)$ the potential function $U(x) = - \int f(x) dx$

- Time Evolution – analytic solution**

$$x(0) = x_0, \quad E = E_0 \quad (\text{initial conditions})$$

$$t = \int_{\chi=x_0}^x \frac{d\chi}{\sqrt{2(E_0 - U(\chi))}} \Rightarrow t = t(x; x_0) \stackrel{\text{inv}}{\Rightarrow} x = x(t; x_0)$$

- Equilibrium solutions (points)** $x(t) = x_0, \quad y(t) = y_0 \quad \forall t$

critical points of vector field

$$y_0 = 0, \quad f(x_0) = 0$$

see [cmath31.nb](#)

$$f(x_0) = 0 \Leftrightarrow f(x) = - \left. \frac{dU}{dx} \right|_{x=x_0} = 0$$

extrema of
potential
function

Part A. Conservative Systems

- **Phase space trajectories** $f(x,y)=c$ are given analytically by the energy integral

$$E = \frac{1}{2} y^2 + U(x) = E_0$$

- Each trajectory corresponds to a constant value of E_0 .
- An energy value E_0 may correspond to more than one trajectory (different orbits with the same energy)
- Phase space trajectories are symmetric with respect to axis $y=0$

$$y = \pm \sqrt{2(E_0 - U(x))}$$

- Flow is directed from left to right for $y>0$ and the opposite for $y<0$
- If a trajectory intersects the axis $y=0$, this intersection should be perpendicular i.e. the vector field on the $y=0$ axis is vertical to it, $\mathbf{f}=(0, f(x))$

Phase space diagram or Phase Portrait : a sufficient collection of trajectories for identifying all possible trajectories

Part A. Conservative Systems

- **Restriction of range of motion** (valid x -interval)

$$E - U(x) \geq 0$$

$-\infty < x < \infty$

unbounded motion to both directions

$x_{\min} \leq x < \infty$

unbounded motion on the right

$-\infty < x \leq x_{\max}$

unbounded motion on the left

$x_{\min} \leq x \leq x_{\max}$

bounded motion

* At $x=x_{\min}$ and $x=x_{\max}$ is $y=0$

boxed text: bounded motion → closed phase space trajectory → periodic motion (oscillation)

- **period of oscillation**

$$T = 2 \int_{x_{\min}}^{x_{\max}} \frac{dx}{\sqrt{2(E - U(x))}}$$

see cmath33.nb

Conservative Linear Systems

1. Harmonic oscillator (elliptic system)

$$\ddot{x} = -kx, \quad k > 0 \quad \text{or} \quad \ddot{x} + \omega^2 x = 0, \quad \omega = \sqrt{k}$$

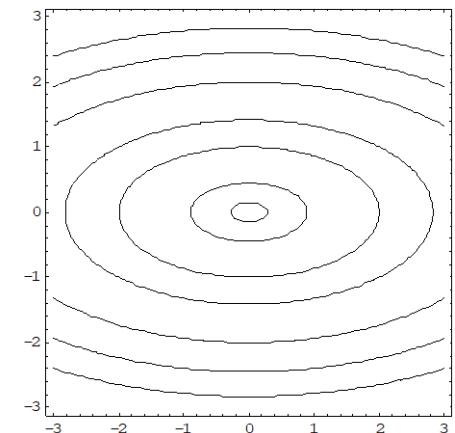
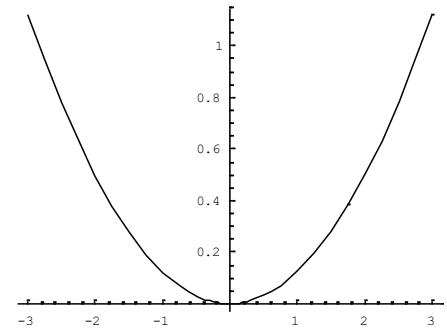
solution : $x(t) = x_0 \cos(\omega t) + \frac{y_0}{\omega} \sin(\omega t), \quad y(t) = y_0 \cos(\omega t) - x_0 \omega \sin(\omega t)$

Energy : $U = \frac{1}{2} \omega^2 x^2, \quad E = \frac{1}{2} y^2 + U,$

trajectories : $\frac{y^2}{2E} + \frac{x^2}{2E/\omega^2} = 1$

Range : $-\frac{\sqrt{2E}}{\omega} \leq x \leq \frac{\sqrt{2E}}{\omega}$

period : $T = \frac{2\pi}{\omega}$ independent of the initial conditions !



Conservative Linear Systems

1. Repulsive force proportional to the distance (hyperbolic system)

$$\ddot{x} = kx, \quad k > 0 \quad \text{or} \quad \ddot{x} - a^2 x = 0$$

solution : $x(t) = x_0 \cosh(at) + \frac{y_0}{a} \sinh(at), \quad y(t) = y_0 \cosh(at) + x_0 a \sinh(at)$

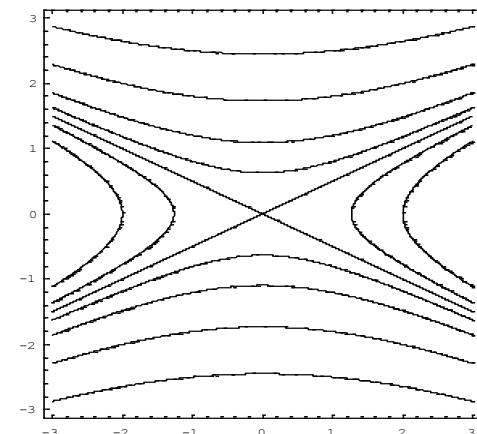
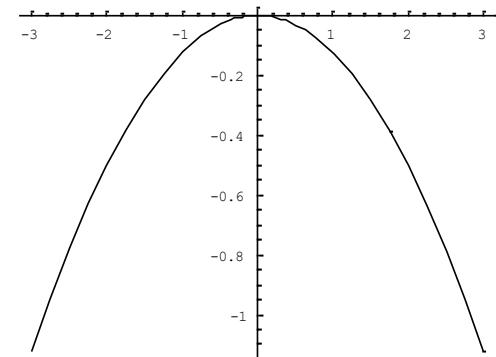
Energy : $U = -\frac{1}{2} a^2 x^2, \quad E = \frac{1}{2} y^2 + U,$

trajectories : $\frac{y^2}{2E} - \frac{x^2}{2E/a^2} = 1$

if $E > 0 \Rightarrow -\infty < x < \infty$

Range : if $E < 0 \Rightarrow x \leq -\frac{\sqrt{-2E}}{a} \quad \text{or} \quad x \geq \frac{\sqrt{-2E}}{a}$

Asymptotic orbits : $E = 0 \Rightarrow y = \pm ax$



Part A. Conservative Systems

Linear stability of equilibrium points

ODEs : $\dot{x} = y, \quad \dot{y} = f(x)$

equilibrium solution (EQP): $x=x_0, y=0$

solution near the equilibrium solution : $x=x_0 + \Delta x, y=0 + \Delta y$

ODEs $\Rightarrow \Delta\dot{x} = \Delta y, \quad \Delta\dot{y} = f(x_0 + \Delta x)$ equations of deviations
from EQP

(after 1st order Taylor expansion of f)

$$\Delta\dot{x} = \Delta y, \quad \Delta\dot{y} = -k\Delta x, \quad k = -\left.\frac{df}{dx}\right|_{x_0} = \frac{d^2V}{dx^2},$$

or $\Delta\ddot{x} = -k\Delta x$

Linear ODE of deviations

k : linear stability index

$k > 0 \Rightarrow$ harmonic oscillations \Rightarrow linearly **stable** EQP

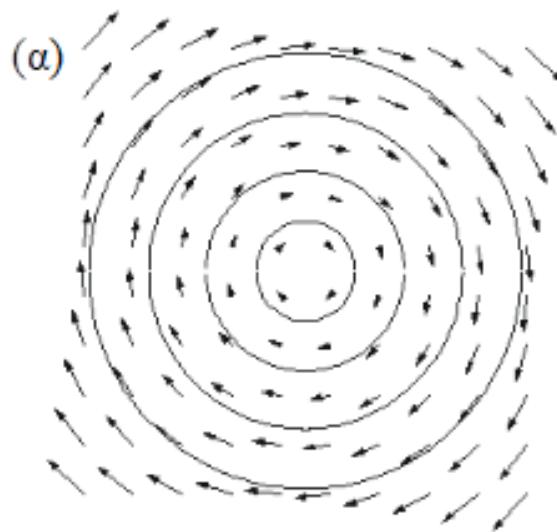
$k < 0 \Rightarrow$ hyperbolic motion \Rightarrow linearly **unstable** EQP

$k = 0 \Rightarrow$ critical case

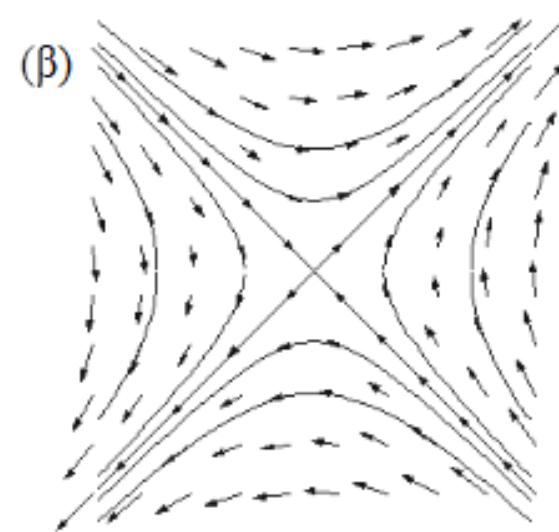
Part A. Conservative Systems

Linear stability of equilibrium points

- Phase space topological structure near Equilibria



Linear stability



Linear instability

Part A. Conservative Systems

General study of Dynamics

- Find EQPs and their linear stability
- Plot The phase space portrait including sufficient orbits (*Energy levels : below, between and above EQPs*).
- Determine the trajectories of asymptotic solutions, *which have the same energy level with the unstable EQPs*. The set of the above asymptotic trajectories (for each EQP) is called the **separatrix curve**
- Separatrix = **stable** + **unstable** manifold
- The separatrix may form **close loops** containing closed-periodic trajectories
- Determine the periods of the periodic trajectories. Close to the stable EQP the period is approximated given by the linear system. As the separatrix is approximated then $T \rightarrow \infty$.

Example 1 : Strongly nonlinear oscillator

[math37example1.nb](#)

Example 2 : 3-mode oscillations in polynomial potential

[math37example2.nb](#)

Example 3 : Oscillations in the Yukawa potential

[math37example3.nb](#)

The dynamics of the simple pendulum

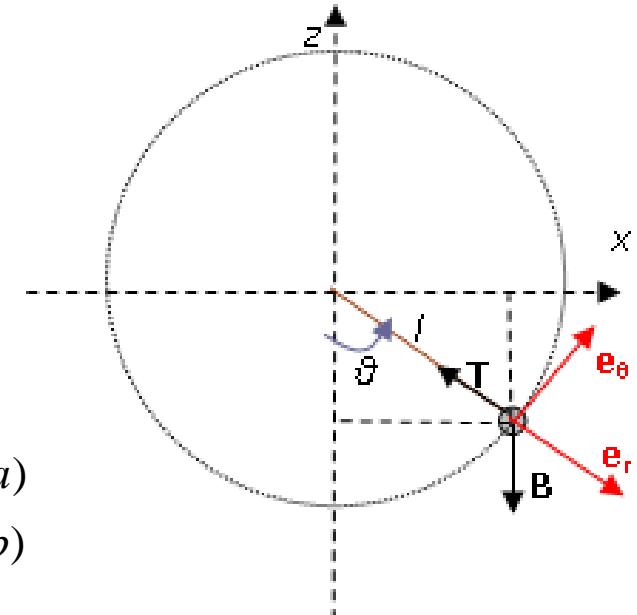
$$m\ddot{\mathbf{r}} = \mathbf{F}, \quad \mathbf{F} = \mathbf{B} + \mathbf{T}, \quad \mathbf{B} = mg\mathbf{k}$$

in polar coordinates

$$\ddot{\mathbf{r}} = (\ddot{r} - r\dot{\theta}^2)\mathbf{e}_r + (2\dot{r}\dot{\theta} + r\ddot{\theta})\mathbf{e}_\theta$$

$$\mathbf{B} = mg \cos \theta \mathbf{e}_r - mg \sin \theta \mathbf{e}_\theta$$

$$\begin{aligned} \mathbf{T} = -T\mathbf{e}_r & \Rightarrow -ml\dot{\theta}^2 = mg \cos \theta - T & (1a) \\ & \Rightarrow ml\ddot{\theta} = -mg \sin \theta & (1b) \end{aligned}$$



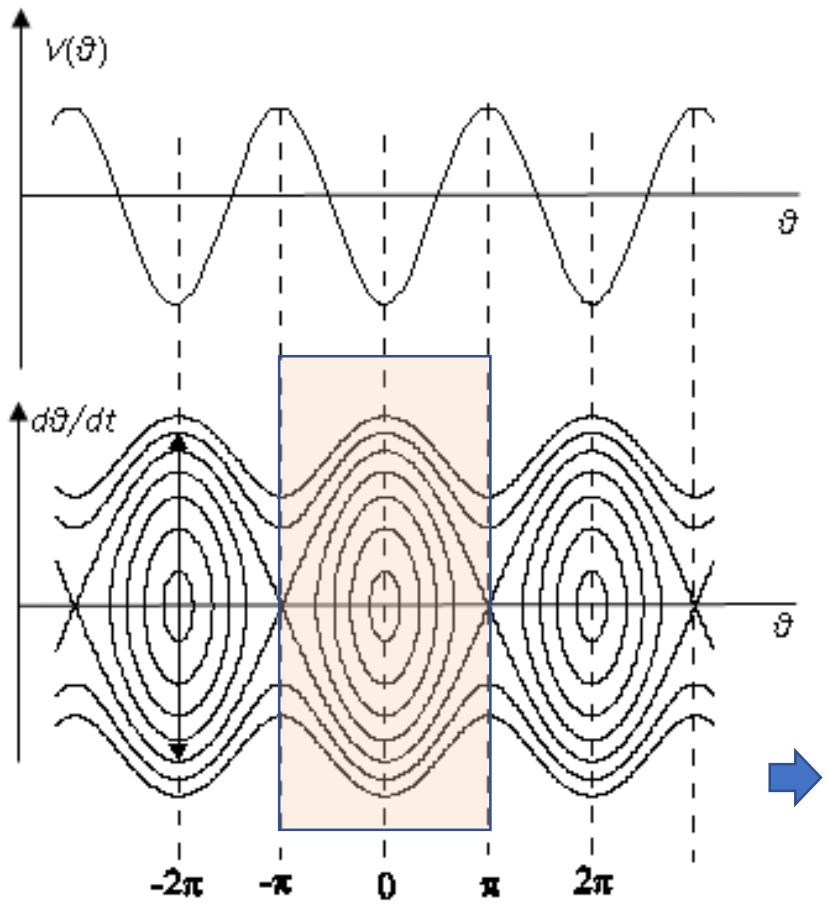
$$T = \frac{1}{2}ml^2\dot{\theta}^2, \quad V = mgz = -mgl \cos \theta \quad \Rightarrow \quad \frac{1}{2}ml^2\dot{\theta}^2 - mgl \cos \theta = E^* \text{ (const)}$$

$$(1b) \Rightarrow \ddot{\theta} = -\omega_0^2 \sin \theta, \quad \omega_0 = \sqrt{g/l}$$

$$f = -\omega_0^2 \sin \theta \Rightarrow V = -\int f d\theta = -\omega_0^2 \cos \theta \Rightarrow$$

$$\frac{1}{2}\dot{\theta}^2 - \omega_0^2 \cos \theta = E \text{ (Energy integral)} \quad (E = E^* / ml^2)$$

The dynamics of the simple pendulum



Equilibrium points : $\sin\theta=0$

$$\theta_0 = 0, 2\pi, 4\pi, \dots$$

$$\theta_0 = \pi, 3\pi, 5\pi, \dots$$

$$V'' = \omega_0^2 > 0$$

$$V'' = -\omega_0^2 < 0$$

ευσταθεια

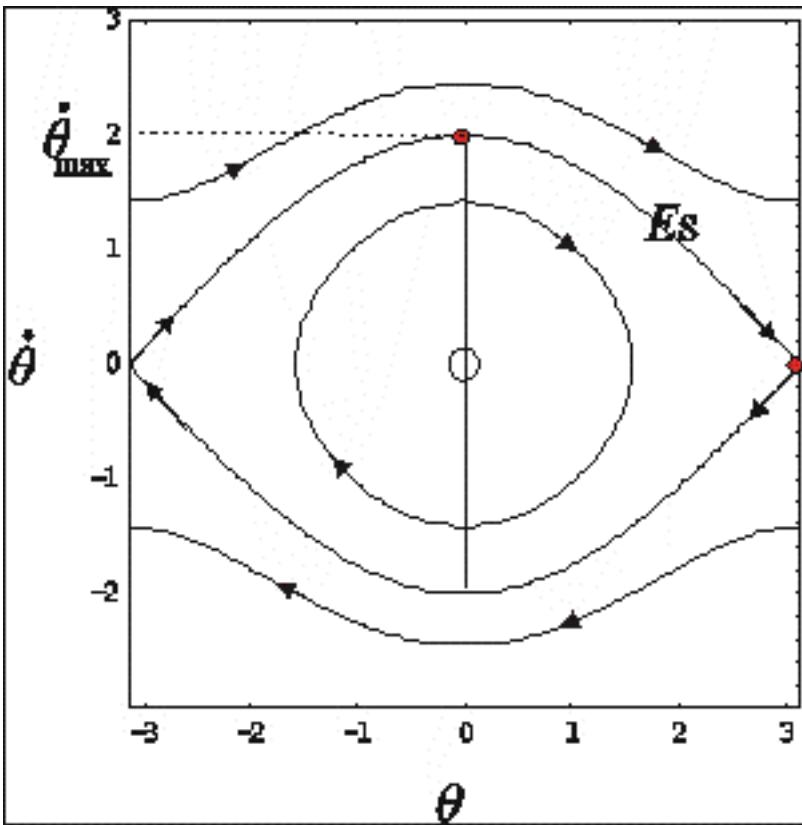
ασταθεια

Remark: We observe that the ODE of motion is invariant under the translation

$$\theta \rightarrow \theta + 2k\pi, \quad k \in \mathbf{Z}$$

i.e. the system can be studied in the interval $-\pi \leq \theta \leq \pi$

The dynamics of the simple pendulum



Energy of Separatrix = Energy of unstable equilibrium point

$$E_s = -\omega_0^2 \cos \pi = \omega_0^2$$

$$\text{Width of Separatrix} \quad \Delta \dot{\theta} = 2\dot{\theta}_{\max} = 4\omega_0$$

Types of motion

$E < E_s$: [Librations](#)

$E > E_s$: [Rotations](#)

$E = E_s$: [Asymptotic motion](#)

Amplitude of libration

$$\theta_{\max} = -\theta_{\min} = \arccos\left(-\frac{E}{\omega_0^2}\right), \quad E < E_s$$

$$A_\theta = \vartheta_{\max} - \vartheta_{\min} = 2\vartheta_{\max}$$

Period of librations

$$T = 4 \int_0^{\theta_{\max}} \frac{d\theta}{\sqrt{2(E + \omega_0^2 \cos \theta)}}$$

Period of Rotation

$$T = 2 \int_0^\pi \frac{d\theta}{\sqrt{2(E + \omega_0^2 \cos \theta)}}$$

Εξάσκηση : Ενδεικτικές Ασκήσεις του Βιβλίου που δεν απαιτούν Mathematica

3.3.1, 3.3.2, 3.3.4 (χωρίς το T), 3.4.1